A Probabilistic Time Reversal Theorem (Poster)

K. Baclawski
Northeastern University, Boston, MA USA

Abstract—Combining independent observations is commonly performed by using a least squares technique, as it is thought that this is necessary to achieve an optimal solution. The purpose of this article is to show that this is not always the case. The particular example combines observations that are exponentially distributed. One application of this technique is to determine the time of a singular event which initiated a set of decay processes having known half-lives. The time of the singular event decays backwards in time with an exponential distribution. We find that the accuracy of this method is significantly better than the accuracy of a least squares technique. The improved accuracy can be important for applications that require combining many noisy observations, such as situation awareness.

Keywords: information fusion, exponential distribution, stochastic processes

I. INTRODUCTION

The combination (or fusion) of independent observations is a fundamental mechanism of probability theory that is important for many purposes. For example, situation awareness (SA) is the perception of environmental elements and events with respect to time or space, the comprehension of their meaning, and the projection into the future [3]. To achieve SA, it is typically necessary to combine (or fuse) multiple, independent sources of noisy information. Because of the noisiness of the sources of information, each observation from a source is a probability distribution that specifies the accuracy of the observation. Fusing the observations requires many issues such as transforming the observations so that they are comparable in time and space and then fusing the probability distributions of the distributions. This article is about the fusion of the distributions.

Since the early nineteenth century when the method of least squares was developed by Gauss [7], it has been a common practice to use a least squares technique to fuse a set of independent observations. While such a method is well-justified and optimal if observations are normally distributed, it is not optimal for all distributions. To illustrate this phenomenon, we consider the case of a set of independent observations that are exponentially distributed, and we show that the additional information about the distributions can have a significant impact on accuracy.

The Gauss-Markov theorem states that in a linear stochastic model whose errors have expectation zero, equal variances and no correlations, the best linear unbiased estimators of the coefficients of the model are the least-squares estimators. [8], [9] This result makes no assumptions about the distributions beyond their means and variances, and so it can only yield simple estimation values, not distributions. When one knows the distributions of the errors, one can not only make an estimate, one can also determine a probability distribution. In other words, both the inputs and the outputs of the information fusion are probability distributions. The normal distribution is the best known example of this process. The distribution of the fusion of a set of independent normally distributed observations is also normally distributed. Moreover, this result does not depend on any other assumptions or results, such as a least squares technique or Bayes’ Law [2], [10]. The fact that the normal distribution is closed under fusion makes it an especially useful distribution for combining independent observations. In this article, it is shown that the case of exponential distributions is also closed under fusion. Whether there are other classes of distributions have this property is an interesting question.

Time reversal is a useful technique for audio and radio signaling and detection applications. [4]–[6] In this technique, a received signal is recorded, reversed, and sent back to the original source. The time reversal theorem considered here differs from the technique for signaling and detection in that we are concerned with stochastic phenomena in which the time of an event (or, more precisely, knowledge of the time of an event) decays exponentially backwards in time. The usefulness of the theorem is due to the improved accuracy of the estimation, not the time reversal itself.

II. THE CLASS OF EXPONENTIAL DISTRIBUTIONS

An exponential distribution models the behavior of a continuous, memoryless waiting time. In other words, if one has been waiting a period of time, then the conditional distribution starting at the end of the period is the same as the distribution starting at the initial time. The behavior is determined by just two constants: the starting point s and the expected duration τ > 0 after the starting point s. Let \( ED(s, \tau) \) denote the distribution in this case. If \( X \) is a random variable whose distribution is \( ED(s, \tau) \), then the probability density of \( X \) is given by:

\[
dens(X = x) = \begin{cases} 
0, & \text{if } x < s, \\
\frac{1}{\tau} e^{-\frac{x-s}{\tau}}, & \text{if } x \geq s.
\end{cases}
\]

The expectation of \( X \) is easily seen to be \( s + \tau \), and the variance is \( \tau^2 \). The exponential distribution models the time between independent occurrences in a sequence that occurs at
a constant rate. The rate is given by $\lambda = \frac{1}{2}$. One can also regard the exponential distribution as modeling a quantity that decays at the rate $\lambda$. When viewed in terms of a decay, one usually takes the base of the exponential to be 2 rather than $e$, and the parameter analogous to the mean duration $\tau$ is the half-life, which is commonly written $t_{1/2}$. The half-life is related to the mean duration by the formula: $t_{1/2} = \ln(2)/\lambda$. The half-life is the median duration, so the median of the distribution is $s + t_{1/2}$.

The constraint that $\tau$ be positive is not essential, and one can easily generalize the class of exponential distributions so that $\tau$ can be any real number. If $\tau = 0$, then the distribution is a discrete distribution with all probability concentrated at $s$ (i.e., the density is a delta function). Having the distribution $ED(s,0)$ means that an observation has value $s$ with probability 1. When $\tau < 0$, the observation decays in the reverse direction. The probability density of a random variable $X$ whose distribution is $ED(s,\tau)$ is given by:

$$dens(X = x) = \begin{cases} \frac{1}{\tau} e^{-\frac{x-s}{\tau}}, & \text{if } x \leq s, \\ 0, & \text{if } x > s. \end{cases}$$

Given a random variable $X$ whose distribution is $ED(s,\tau)$, the random variable $X - s$ will have distribution $ED(0,\tau)$, and the mean of $X - s$ is $\tau$. If $\tau$ is nonzero, then $X - s$ has distribution $ED(0,1)$. As already discussed above, the median of $ED(s,\tau)$ is $s + \ln(2)/\tau$. Because this distribution is so highly asymmetric, it is more reasonable to use the median rather than the mean when it is necessary to reduce the distribution to an estimate consisting of a single number.

III. INFORMATION FUSION OF EXPONENTIAL DISTRIBUTIONS

We now compute the distribution of the fusion of a set of $N$ independent exponentially distributed random variables. The following is the general case:

**Theorem 3.1:** Let $X_i$ be a set of independent random variables whose distributions are $ED(s_i, \tau_i)$, for $i = 1, \ldots, N$. Write $Y$ for the fusion of the random variables $\{X_i\}$.

1. If $\tau_i > 0$ for every $i$, then $Y$ has the distribution $ED(\max_{i=1}^N (s_i), \left(\sum_1^N (\frac{1}{\tau_i})^{-1}\right))$.
2. If $\tau_i < 0$ for every $i$, then $Y$ has the distribution $ED(\min_{i=1}^N (s_i), \left(\sum_1^N (\frac{1}{\tau_i})^{-1}\right))$.

**Proof.** We first consider the case in which all the $\tau_i$ are positive. The probability density of $Y$ is given by normalizing the product of the densities of the random variables $X_i$ as follows (see [1]):

$$dens(Y = y) = \begin{cases} 0, & \text{if } y < s_i \text{ for some } i, \\ C\prod_{i=1}^N e^{-\frac{y-s_i}{\tau_i}}, & \text{if } y \geq s_i \text{ for all } i, \end{cases}$$

where $C$ is the normalization constant such that this formula defines a probability density function. Now $y \geq s_i$ for all $i$ if and only if $y \geq \max_i(s_i)$. Let $s = \max_i(s_i)$. The formula above may then be written as follows:

$$dens(Y = y) = \begin{cases} 0, & \text{if } y < s, \\ Ce^{-\sum_{i=1}^N (\frac{y-s_i}{\tau_i})}, & \text{if } y \geq s. \end{cases}$$

The exponent in the formula above may be written as follows:

$$-\sum_{i=1}^N \left(\frac{y-s_i}{\tau_i}\right) = -\sum_{i=1}^N \left(\frac{y}{\tau_i} - \frac{s_i}{\tau_i}\right)$$

$$= -\sum_{i=1}^N \left(\frac{y}{\tau_i}\right) + \sum_{i=1}^N \left(\frac{s_i}{\tau_i}\right)$$

$$= -y\sum_{i=1}^N \left(\frac{1}{\tau_i}\right) + \sum_{i=1}^N \left(\frac{s_i}{\tau_i}\right)$$

Let

$$\tau = \frac{1}{\sum (\frac{1}{\tau_i})},$$

or in terms of the rate parameters, $\lambda = \sum_{i=1}^N \lambda_i$. Also let

$$D = -\frac{s}{\tau} + \sum_{i=1}^N \left(\frac{s_i}{\tau_i}\right).$$

Note that $D$ does not depend on $y$. One can then write the sum above as follows:

$$-\sum_{i=1}^N \left(\frac{y-s_i}{\tau_i}\right) = -y\tau + \sum_{i=1}^N \left(\frac{s_i}{\tau_i}\right)$$

$$= -y + \frac{s}{\tau} - \frac{s}{\tau} + \sum_{i=1}^N \left(\frac{s_i}{\tau_i}\right)$$

$$= -y + s + D$$

The probability density when $y \geq s$ may then be computed as follows:

$$Ce^{-\sum_{i=1}^N \left(\frac{y-s_i}{\tau_i}\right)} = Ce^{-\left(\frac{y-s}{\tau}\right) + D}$$

$$= Ce^{D} e^{-\left(\frac{y-s}{\tau}\right)}$$

Since $Ce^D$ is a constant independent of $y$, it follows that $Y$ is exponentially distributed with distribution $ED(s, \tau)$. In the same way, one can also compute the distribution of the fusion when all of the $\tau_i$ are negative. The only difference is that the maximum of the $\{s_i\}$ is replaced by the minimum. The result then follows.
In the special case where all of the \( \tau_i \) are the same, the formula simplifies to the following:

**Corollary 3.2:** Let \( X_i \) be a set of independent random variables whose distributions are \( ED(s_i, \tau) \), for \( i = 1, \ldots, N \). Write \( Y \) for the fusion of the random variables \( \{X_i\} \).

1. If \( \tau > 0 \), then \( Y \) has the distribution \( ED(\max_{i=1}^{N} (s_i, \tau)) \).
2. If \( \tau < 0 \), then \( Y \) has the distribution \( ED(\min_{i=1}^{N} (s_i, \tau)) \).

The case in which some of the \( \tau_i \) are positive and some are negative does not result in an exponential distribution, so it is not considered.

**IV. Experimental Observations**

We now consider an experiment in which a particle splits into a set of \( N \) particles, each of which subsequently decays independently. Suppose that the half-lives of the products are known to be \( \{h_i|1 \leq i \leq N\} \) and that one can observe the time when each of the products decays. The splitting of the original particle is called the singular event. The problem is to determine the time when the singular event occurred.

**Theorem 4.1:** If a singular event results in \( N \) independent exponential decay processes with half lives \( \{h_i|1 \leq i \leq N\} \), then the time of the singular event decays backwards in time from the time when the first decay product is observed, with a half-life equal to \( (\sum(h_i^{-1}))^{-1} \).

**Proof.** Let \( t \) be the time when the original particle split, and let \( \{t_i\} \) be the times when the product particles are observed to decay. Then the duration from \( t \) to \( t_i \) has the distribution \( ED(t, h_i/\ln(2)) \), and the duration from \( t \) to \( t_i \) has the distribution \( ED(t_i, -h_i/\ln(2)) \). The latter set of distributions are all observing the same time \( t \), and they are assumed to be independent, so Theorem 3.1 applies. Therefore, the distribution of the time \( t \) is

\[
ED\left( \min_{i=1}^{N} (t_i), \frac{1}{\sum(h_i^{-1})} \right) = ED\left( \min_{i=1}^{N} (t_i), \frac{1}{\sum(h_i^{-1})} \right)/\ln(2).
\]

The result then follows.

In the special case where all of the half lives are the same, the formula simplifies as follows:

**Corollary 4.2:** If a singular event results in \( N \) independent identically distributed exponential decay processes with half-life \( t_{1/2} \), then the time of the singular event decays backwards in time from the time when the first decay product is observed, with a half-life equal to \( t_{1/2}/N \).

**V. The Least Squares Method**

We now apply the least squares method to the scenario of Section IV above. For simplicity we consider only the case where the \( N \) exponential distributions have the same half-life \( t_{1/2} \). The mean lifetime is \( \tau = t_{1/2}/\ln(2) \), and the variance is \( \tau^2 \). Let \( x_i \) be the actual decay observation for the \( i^{th} \) product particle. Each decay observation \( x_i \) may be written in the form \( t + \tau + e_i \), where \( e_i \) is a random “error” having mean 0 and variance \( \tau^2 \). Therefore each \( X_i - \tau \) has mean \( t \) and variance \( \tau^2 \). The least squares estimate for \( t \) is the sample average \( \frac{1}{N} \sum_{i=1}^{N} (x_i - \tau) = \left( \frac{1}{N} \sum_{i=1}^{N} x_i \right) - \tau \). This estimate has variance \( \tau^2/N \) and standard deviation \( \tau/\sqrt{N} \). The information fusion and the least squares method are compared in Figure 2.

<table>
<thead>
<tr>
<th></th>
<th>Least Squares</th>
<th>Information Fusion</th>
</tr>
</thead>
<tbody>
<tr>
<td>Estimate</td>
<td>( \frac{1}{N} \sum_{i=1}^{N} x_i - \tau )</td>
<td>( \min_{i=1}^{N} (x_i) - \frac{\ln(2)\tau}{N} )</td>
</tr>
<tr>
<td>Variance</td>
<td>( \tau^2/N )</td>
<td>( \tau^2/N^2 )</td>
</tr>
<tr>
<td>Std Dev</td>
<td>( \tau/\sqrt{N} )</td>
<td>( \tau/N )</td>
</tr>
</tbody>
</table>

![Fig. 1. Information Fusion versus the Least Squares Method for N=4](image)

![Fig. 2. Comparison of Information Fusion and Least Squares Methods](image)

It is apparent that information fusion will be more accurate than the least squares technique, especially for large \( N \). For example, consider the case in which \( t = 5 \), \( \tau = 1 \) and \( N = 4 \). Information fusion yields the distribution \( ED(\min(t_i), -0.25) \). The median value for \( \min(t_i) \) is \( t + \ln(2)/4 = 5.17329 \). If this experiment is run many times, the values for \( \min(t_i) \) will vary. To show a representative sample of the behavior, the quartiles are shown for the two methods in Figure 1, and in the later graphs. The peaks of the exponential distribution should all be the same, but the graphs do not show this due to round-off errors.

The least squares method yields the estimate \( \left( \frac{1}{N} \sum t_i \right) - 1 \). This estimate is normally distributed with mean 5 and standard deviation 0.5. Accordingly it will be within 0.3372 of 5 about
The quartiles are therefore at 4.6628, 5 and 5.3372. These are shown in Figure 1. The information fusion distribution is approximately 4 times more accurate than the least squares distribution. Half of the improvement is due to the fact that the normal distribution is a two-tailed distribution while the exponential distribution is a single-tailed distribution. The other half of the improvement is the result of the improvement in the variance.

In Figure 3 the density functions are compared for the case in which $N = 16$. In this case, the information fusion distribution is approximately 8 times more accurate than the least squares distribution. In general the information fusion distribution will be approximately $2\sqrt{N}$ times more accurate. The factor of 2 in this approximation is somewhat arbitrary, and represents the fact that the exponential distribution is single-tailed, while the normal distribution is double-tailed. Based on the estimates alone, one might take the factor to be 1.4. See Figure 5 which compares the estimates for various values of $N$.

In Figure 4 the density functions are compared for the case in which $N = 1$. Here there is no information fusion, so it is only comparing the exponential distribution with the normal distribution.

VI. CONCLUSION

We have shown an example for which one can significantly improve the accuracy of the fusion of observations when their probability distributions are exponential, compared with techniques, such as a least squares technique, that do not require specific knowledge of the probability distributions. Using information fusion of the distributions rather than a least squares estimation has the further advantage that the result of the method is a probability distribution rather than an estimated value and a standard deviation that bounds the error. Consequently, the result may, in principle, be used as input for further computations that also use fuse the full probability distributions. Combining several such operations will ultimately enable improved results for applications that require such complex analyses, such as situation awareness.

The fusion of exponential distributions and least squares methods were compared for an example in which a singular event initiates an independent set of decay processes having known half-lives. The result can be viewed as reversing time, because our knowledge of the time of the singular event decays backwards in time with an exponential distribution whose parameters were computed. The accuracy of this technique was found to be significantly better than the accuracy obtained by using a least squares method.

Fusion of precisely known distributions has many advantages, but there are only a few examples of classes of probability distributions to which the technique is currently being applied. Introducing additional classes of distributions that are
closed under fusion could have a significant impact on the accuracy of estimation in a variety of domains.

VII. ACKNOWLEDGMENTS

Initial research that led to this paper was performed while the author was on sabbatical at the Division of Preventive Medicine, Brigham and Women’s Hospital; and the author wishes to thank Tianhua Niu, JoAnn Manson and Simin Liu for their help and encouragement. I also appreciate the encouragement and support of colleagues at Versatile Information Systems and Composable Logic, including Mitch Kokar and Jeff Smith. I especially appreciate Jeff Keller for carefully proofreading the manuscript.

REFERENCES